

An indefinite integral equation without irregular frequencies for the floating-body problem

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In his classic work [1] on the floating-body problem (FBP) for water of constant finite depth d , John reduced this boundary-value problem (BVP) to an integral equation (IE) on the wetted portion S of body's surface. For this purpose he constructed Green's function $G(P, Q)$, $P = (x, y, z)$, $Q = (\xi, \eta, \zeta)$, such that: (i) $G - |P - Q|^{-1}$ satisfies the Laplace equation in the water layer $L = \{-\infty < x, z < +\infty, -d < y < 0\}$; (ii) the radiation condition holds for G as $|P - Q| \rightarrow \infty$; (iii) the following boundary conditions are fulfilled: $[G_y - (\omega^2/g)G]_{y=0} = 0$ and $G_y|_{y=-d} = 0$, where ω is the radian time-frequency of waves and g is the acceleration due to gravity. Then the single-layer potential (SLP) – see the first term in formula (1) below – satisfies all conditions of the FBP with exception of the non-homogeneous Neumann condition (NC) on S . In order to satisfy this NC one applies the jump property of the normal derivative of the SLP and arrives at the Fredholm IE for μ , provided S has the following properties: (a) S has no common points with $\{y = -d\}$; (b) S is attached to $\{y = 0\}$ forming a C^2 -surface with the mirror image of S in $\{y = 0\}$. However, thus obtained IE has the following drawback discovered by John himself. At a certain sequence of frequencies (these *irregular* frequencies are related to the spectrum of a BVP in the domain confined between $\{y = 0\}$ and S) the IE is not solvable for all right-hand-side terms (RHST) even though the FBP is uniquely solvable, which was proved by John for S satisfying the following additional condition: (c) S lies within the vertical cylinder whose generators go through the water-line along which S is attached to the plane $\{y = 0\}$.

Since the publication of John's paper, the question of formulating IE uniquely solvable for all frequencies has occupied one of the central positions in the studies of the FBP (see [2], Sections 3.1.1 and 3.1.2). The aim of the present work is to formulate a new, free of irregular frequencies IE for the FBP. This IE belongs to the class of indefinite equations introduced by S. G. Krein in the 1950's and considered in his book [3].

Assuming that S satisfies conditions (a)–(c) and the RHST f in the NC on S belongs to $C^{0,\alpha}(S)$, let us seek a solution to the FBP in the form:

$$\int_S \mu(Q) G(P, Q) dS_Q + \int_S \nu(Q) \frac{\partial G(P, Q)}{\partial n_Q} dS_Q, \quad (1)$$

where n_Q is the unit normal pointing into the water domain at $Q \in S$, and $\mu \in C^{0,\alpha}(S)$ and $\nu \in C^{1,\alpha}(S)$ are unknown densities. Then the normal derivative of the function given by formula (1) exists for $P \in S$, and one arrives at the required IE:

$$-\mu(P) + \frac{1}{2\pi} \int_S \mu(Q) \frac{\partial G(P, Q)}{\partial n_P} dS_Q + \frac{1}{2\pi} \frac{\partial}{\partial n_P} \int_S \nu(Q) \frac{\partial G(P, Q)}{\partial n_Q} dS_Q = f(P) \quad \text{or} \quad (K - I)\mu + T\nu = f \quad (2)$$

for short. Here I is the identity operator, K and its adjoint K' are compact operators in $C(S)$, and T is self-adjoint when defined in the manner described in [4], Section 2.7. According to results of Section 19 in [3], equation (2) is solvable for all RHST, if the intersection of the null-spaces of $K' - I$ and T is the zero subspace. In order to prove this let us suppose that μ_0 satisfies equations $(K' - I)\mu_0 = 0$ and $T\mu_0 = 0$ simultaneously. Then the first of these equations implies that $\mu_0 \in C^{1,\alpha}(S)$, and so the second equation means that the double-layer potential in (1) with $\nu = \mu_0$ solves the homogeneous FBP. By virtue of the uniqueness theorem for the FBP, this potential is identically equal to zero in the water domain. Then applying the jump formula as P tends to its limit position on S , one gets that $(K' + I)\mu_0 = 0$, which combined with $(K' - I)\mu_0 = 0$ gives that $\mu_0 = 0$ on S . Hence the following result is true.

Let conditions (a)–(c) hold. Then the indefinite equation (2) is solvable for all RHST. Substituting a solution of the IE (2) into the representation formula (1), one obtains a unique solution of the FBP.

There are two subspaces such that if a solution (μ, ν) belongs to their direct product, then this solution is unique. The description of these subspaces can be given in terms of ranges and null-spaces of $K' - I$ and T .

- [1] F. John, On the motion of floating bodies. II, *Comm. Pure Appl. Math.* **3** (1950) 45–101.
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- [4] D. Colton, R. Kress, *Integral Equation Methods in Scattering Theory*, Wiley, 1983.